

Κβαντομηχανική II, ΣΕΜΦΕ

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Άσκηση 1.

α)

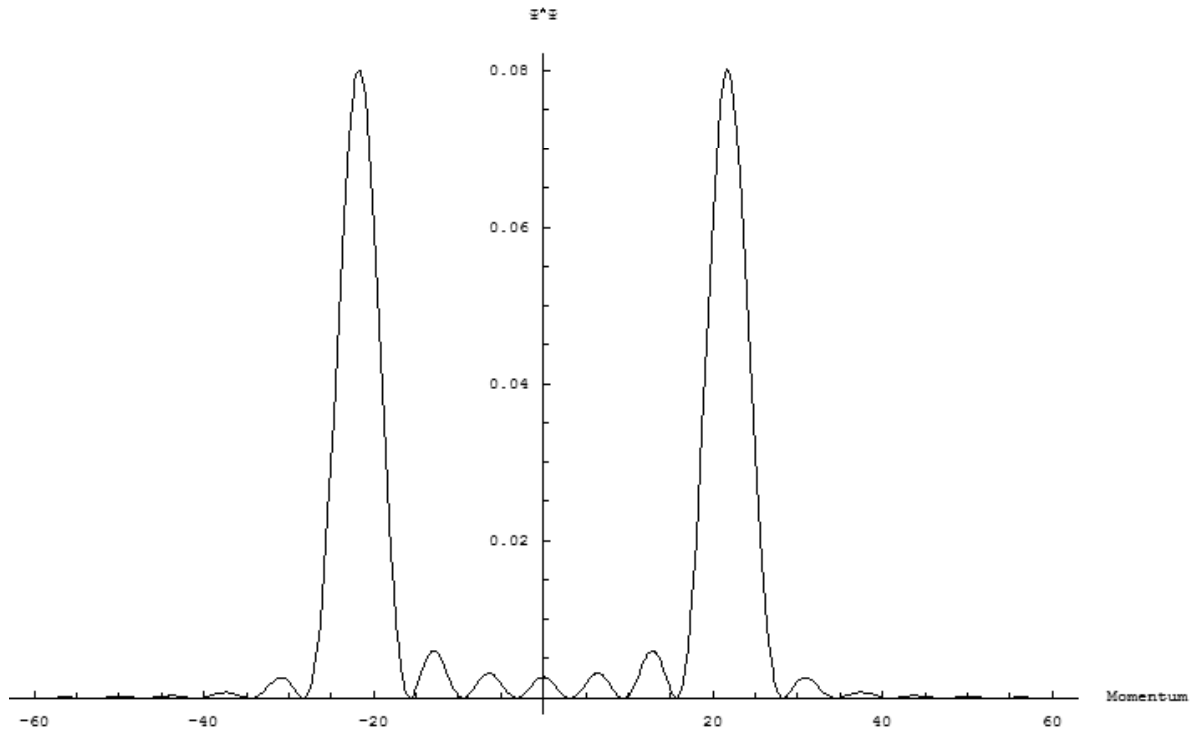
$$\begin{aligned} N^2 \int_0^L \sin^2\left(\frac{n\pi x}{L}\right) dx &= 1 \\ \int_0^L \sin^2\left(\frac{n\pi x}{L}\right) dx &= \int_0^L \frac{1}{2} dx - \frac{1}{2} \int_0^L \cos\left(\frac{2n\pi x}{L}\right) dx = \frac{L}{2} \\ \Rightarrow N^2 \frac{L}{2} &= 1 \Rightarrow N = \sqrt{\frac{2}{L}} \quad \blacksquare \end{aligned}$$

β)

$$\begin{aligned} \Psi(x) &= \frac{i}{2} \sqrt{\frac{2}{L}} \left(e^{-i\frac{n\pi x}{L}} - e^{i\frac{n\pi x}{L}} \right) \\ \Phi(p) &= \frac{-i}{\sqrt{2\pi\hbar}} \frac{1}{\sqrt{2L}} \int_0^L e^{-i\frac{px}{\hbar}} e^{i\frac{n\pi x}{L}} dx + \frac{i}{\sqrt{2\pi\hbar}} \frac{1}{\sqrt{2L}} \int_0^L e^{-i\frac{px}{\hbar}} e^{-i\frac{n\pi x}{L}} dx = \\ &= \frac{-i}{\sqrt{4\pi\hbar L}} \int_0^L e^{i\left(\frac{n\pi}{L} - \frac{p}{\hbar}\right)x} dx + \frac{i}{\sqrt{4\pi\hbar L}} \int_0^L e^{-i\left(\frac{n\pi}{L} + \frac{p}{\hbar}\right)x} dx = \\ &= \frac{-i}{\sqrt{4\pi\hbar L}} \frac{e^{i\left(\frac{n\pi}{L} - \frac{p}{\hbar}\right)L} - 1}{i\left(\frac{n\pi}{L} - \frac{p}{\hbar}\right)} + \frac{i}{\sqrt{4\pi\hbar L}} \frac{e^{-i\left(\frac{n\pi}{L} + \frac{p}{\hbar}\right)L} - 1}{-i\left(\frac{n\pi}{L} + \frac{p}{\hbar}\right)} = \\ &= \frac{-1}{\sqrt{4\pi\hbar L}} \frac{e^{-i\frac{pL}{\hbar}} - 1}{\left(\frac{n\pi}{L} - \frac{p}{\hbar}\right)} - \frac{1}{\sqrt{4\pi\hbar L}} \frac{e^{-i\frac{pL}{\hbar}} - 1}{\left(\frac{n\pi}{L} + \frac{p}{\hbar}\right)} = \\ &= \frac{-1}{\sqrt{4\pi\hbar L}} \left(e^{-i\frac{pL}{\hbar}} - 1 \right) \frac{2\frac{n\pi}{L}}{\left(\left(\frac{n\pi}{L}\right)^2 - \left(\frac{p}{\hbar}\right)^2\right)} \quad \blacksquare \end{aligned}$$

γ)

$$\begin{aligned} |\Phi|^2 &= \Phi^*(p)\Phi(p) = \frac{4\frac{n^2\pi^2}{L^2}}{4\pi\hbar L} \frac{\left(e^{i\frac{pL}{\hbar}} - 1\right)\left(e^{-i\frac{pL}{\hbar}} - 1\right)}{\left(\left(\frac{n\pi}{L}\right)^2 - \left(\frac{p}{\hbar}\right)^2\right)^2} = \\ &= \frac{4n^2\pi}{\hbar L^3} \frac{\sin^2\left(\frac{pL}{2\hbar}\right)}{\left(\frac{n\pi}{L} - \frac{p}{\hbar}\right)^2 \left(\frac{n\pi}{L} + \frac{p}{\hbar}\right)^2} \quad \blacksquare \end{aligned}$$



Φιγυρε 1: Η γραφική παράσταση της πυκνότητας πιθανότητας συναρτήσεως της ορμής για $\hbar = 1$, $L = 1$ και $n = 10$

Άσκηση 2.

α)

$$\begin{aligned}
 \bullet \quad \langle x \rangle e &= \int_{-\infty}^{+\infty} x \Psi^*(x) \Psi(x) dx = \\
 &= \sqrt{\frac{a}{\pi}} \int_{-\infty}^{+\infty} x e^{-a(x-x_0)^2} dx = \\
 &= \sqrt{\frac{a}{\pi}} \int_{-\infty}^{+\infty} (y + x_0) e^{-ay^2} dy = \\
 &= 0 + x_0 \sqrt{\frac{a}{\pi}} \int_{-\infty}^{+\infty} e^{-ay^2} dy = x_0 \sqrt{\frac{a}{\pi}} \sqrt{\frac{\pi}{a}} = x_0 \\
 \bullet \quad \langle x^2 \rangle &= \sqrt{\frac{a}{\pi}} \int_{-\infty}^{+\infty} x^2 e^{-a(x-x_0)^2} dx = \sqrt{\frac{a}{\pi}} \int_{-\infty}^{+\infty} (y + x_0)^2 e^{-ay^2} dy = \\
 &= \sqrt{\frac{a}{\pi}} \left[\int_{-\infty}^{+\infty} y^2 e^{-ay^2} dy + 2x_0 \int_{-\infty}^{+\infty} y e^{-ay^2} dy + x_0^2 \int_{-\infty}^{+\infty} e^{-ay^2} dy \right] = \\
 &= \sqrt{\frac{a}{\pi}} \left[\frac{1}{2a} \sqrt{\frac{\pi}{a}} + 0 + x_0^2 \sqrt{\frac{\pi}{a}} \right] = x_0^2 + \frac{1}{2a} \\
 \bullet \quad \langle p \rangle &= \int_{-\infty}^{+\infty} \Phi(x) e^{-i\frac{p_0 x}{\hbar}} (-i\hbar) \frac{d\Phi}{dx} dx = \\
 &= \int_{-\infty}^{+\infty} \Phi(x) e^{-i\frac{p_0 x}{\hbar}} (-i\hbar) \left[\frac{d\Phi}{dx} e^{i\frac{p_0 x}{\hbar}} + \left(\frac{ip_0}{\hbar} \right) \Phi e^{i\frac{p_0 x}{\hbar}} \right] dx =
 \end{aligned}$$

$$\begin{aligned}
&= (-i\hbar) \int_{-\infty}^{+\infty} \Phi(x) \frac{d\Phi}{dx} dx + p_0 \int_{-\infty}^{+\infty} \Phi^2 dx = \\
&= \frac{-i\hbar}{2} \int_{-\infty}^{+\infty} \frac{d\Phi^2}{dx} dx + p_0 = \frac{-i\hbar}{2} \Phi^2 \Big|_{-\infty}^{+\infty} + p_0 = 0 + p_0 = p_0
\end{aligned}$$

- $\langle p^2 \rangle = -\hbar^2 \int_{-\infty}^{+\infty} \Psi^*(x) \Psi''(x) dx =$
- $= -\hbar^2 \int_{-\infty}^{+\infty} \frac{d}{dx} (\Psi^* \Psi') dx + \hbar^2 \int_{-\infty}^{+\infty} \Psi'^* \Psi' dx =$
- $= -\hbar^2 \Psi^*(x) \Psi'(x) \Big|_{-\infty}^{+\infty} + \hbar^2 \int_{-\infty}^{+\infty} \Psi'^* \Psi' dx$
- $= \hbar^2 \int_{-\infty}^{+\infty} \Psi'^* \Psi' dx$
- \Rightarrow
- $\langle p^2 \rangle = \hbar^2 \int_{-\infty}^{+\infty} \left[\Phi'(x) e^{-i\frac{p_0 x}{\hbar}} - \frac{ip_0}{\hbar} \Phi(x) e^{-i\frac{p_0 x}{\hbar}} \right] \left[\Phi'(x) e^{i\frac{p_0 x}{\hbar}} + \frac{ip_0}{\hbar} \Phi(x) e^{i\frac{p_0 x}{\hbar}} \right] dx =$
- $= \hbar^2 \int_{-\infty}^{+\infty} (\Phi')^2 dx - ip_0 \hbar \int_{-\infty}^{+\infty} \Phi \Phi' dx + ip_0 \hbar \int_{-\infty}^{+\infty} \Phi' \Phi dx + p_0^2 \int_{-\infty}^{+\infty} \Phi^2 dx =$
- $= p_0^2 + \hbar^2 \sqrt{\frac{a}{\pi}} \int_{-\infty}^{+\infty} a^2 (x - x_0)^2 e^{-a(x-x_0)^2} dx =$
- $= p_0^2 + a^2 \hbar^2 \left(x_0^2 + \frac{1}{2a} - 2x_0^2 + x_0^2 \right) =$
- $= p_0^2 + \hbar^2 \frac{a}{2} \quad \blacksquare$

$\beta)$

- $(\Delta x)^2 = \langle x^2 \rangle - \langle x \rangle^2 = x_0^2 + \frac{1}{2a} - x_0^2 = \frac{1}{2a}$
- $(\Delta p)^2 = \langle p^2 \rangle - \langle p \rangle^2 = p_0^2 + \hbar^2 \frac{a}{2} - p_0^2 = \hbar^2 \frac{a}{2}$
- $(\Delta x) (\Delta p) = \frac{1}{\sqrt{2a}} \sqrt{\frac{a}{2}} \hbar = \frac{\hbar}{2} \quad \blacksquare$

Άσκηση 3.

Ορίζουμε πρώτα την κυματοσυνάρτηση στον χώρο των ορμών

$$\begin{aligned}\Phi(p) &= \int_{-\infty}^{+\infty} \Psi_p^*(x) \Psi(x, 0) dx = \frac{1}{\sqrt{2\pi\hbar}} \left(\frac{\lambda}{\pi}\right)^{\frac{1}{4}} \int_{-\infty}^{+\infty} e^{-\lambda \frac{x^2}{2}} e^{-i \frac{px}{\hbar}} dx = \\ &= \frac{1}{\sqrt{2\pi\hbar}} \left(\frac{\lambda}{\pi}\right)^{\frac{1}{4}} \sqrt{\frac{2\pi}{\lambda}} e^{-\frac{p^2}{2\lambda\hbar^2}} = \left(\frac{1}{\lambda\pi\hbar^2}\right)^{\frac{1}{4}} e^{-\frac{p^2}{2\lambda\hbar^2}}\end{aligned}$$

Έχουμε:

$$\int_{-\infty}^{+\infty} e^{-\alpha x^2 + \beta x} dx = \sqrt{\frac{\pi}{\alpha}} e^{\frac{\beta^2}{4\alpha}}$$

με $\text{Re}(\alpha) > 0$ και $\beta = \text{οτιδήποτε}$

Ορίζουμε την $\Psi(x, t)$

$$\begin{aligned}\Psi(x, t) &= \int_{-\infty}^{+\infty} \Phi(p) \Psi_p(x) e^{-i \frac{E_p t}{\hbar}} dp = \\ &= \left(\frac{1}{\lambda\pi\hbar^2}\right)^{\frac{1}{4}} \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{+\infty} e^{-\frac{p^2}{2\lambda\hbar^2}} e^{i \frac{px}{\hbar}} e^{-i \frac{p^2 t}{2m\hbar}} dp = \\ &= \frac{1}{\sqrt{2\pi\hbar}} \frac{1}{(\lambda\pi\hbar^2)^{\frac{1}{4}}} \int_{-\infty}^{+\infty} e^{-\frac{1}{2\lambda\hbar^2} + \frac{it}{2m\hbar}} p^2 + \frac{ix}{\hbar} p dp = \\ &= \frac{1}{\sqrt{2\pi\hbar}} \frac{1}{(\lambda\pi\hbar^2)^{\frac{1}{4}}} \sqrt{\frac{\pi}{\left(\frac{1}{2\lambda\hbar^2} + \frac{it}{2m\hbar}\right)}} e^{\left(\frac{ix}{\hbar}\right)^2 \frac{1}{4\left(\frac{1}{2\lambda\hbar^2} + \frac{it}{2m\hbar}\right)}} = \\ &= N(t) e^{-\frac{\lambda(t)}{2} x^2}\end{aligned}$$

όπου

$$N(t) = \frac{1}{\sqrt{2\pi\hbar}} \frac{1}{(\lambda\pi\hbar^2)^{\frac{1}{4}}} \sqrt{\frac{\pi}{\left(\frac{1}{2\lambda\hbar^2} + \frac{it}{2m\hbar}\right)}}$$

και

$$\begin{aligned}\lambda(t) &= \frac{1}{2\hbar^2} \frac{1}{\frac{1}{2\lambda\hbar^2} + \frac{it}{2m\hbar}} = \frac{1}{2\hbar^2} \frac{4\lambda m \hbar^3}{2m\hbar + i2\lambda t \hbar^2} \\ &= \frac{\lambda m}{m + i\lambda t \hbar} = \frac{\lambda m(m - i\lambda t \hbar)}{m^2 + \lambda^2 t^2 \hbar^2} = \lambda_1(t) - i\lambda_2(t)\end{aligned}$$

με

$$\lambda_1(t) = \frac{\lambda m^2}{m^2 + \lambda^2 t^2 \hbar^2}$$
$$\lambda_2(t) = \frac{\lambda^2 t \hbar m}{m^2 + \lambda^2 t^2 \hbar^2}$$

Άρα

$$\Psi(x, t) = N(t) e^{i \frac{\lambda_2(t)}{2} x^2} e^{-\frac{\lambda_1(t)}{2} x^2}$$

όπου $\lambda_1(t) \rightarrow 0$ για $t \rightarrow \infty$

Επίσης

$$P(x, t) = |\Psi(x, t)|^2 = |N(t)|^2 e^{-\lambda_1 x^2}$$
$$|N(t)|^2 = \sqrt{\frac{\lambda_1(t)}{\pi}}$$

και

$$(\Delta x)_t = \frac{1}{\sqrt{2\lambda_1(t)}} = \frac{1}{\sqrt{2\lambda}} \sqrt{1 + \frac{\lambda^2 \hbar^2 t^2}{m^2}} \xrightarrow{t \rightarrow \infty} \infty$$
$$(\Delta x)_{t=0} = \frac{1}{\sqrt{2\lambda}}$$

Άρα το κυματοπακέτο απλώνει με τον χρόνο και η πιθανότητα $P(x, t) \rightarrow 0$ για $t \rightarrow \infty$. ■

Άσκηση 4.

$$-\frac{\hbar^2}{2m} \nabla^2 \Psi + V_1 \Psi + iV_2 \Psi = i\hbar \frac{\partial \Psi}{\partial t}$$
$$-\frac{\hbar^2}{2m} \nabla^2 \Psi^* + V_1 \Psi^* - iV_2 \Psi^* = -i\hbar \frac{\partial \Psi^*}{\partial t}$$

$$\frac{\partial(\Psi^* \Psi)}{\partial t} = \frac{\partial \Psi^*}{\partial t} \Psi + \Psi^* \frac{\partial \Psi}{\partial t} =$$
$$= \frac{\hbar}{2im} \Psi \nabla^2 \Psi^* - \frac{1}{i\hbar} V_1 \Psi^* \Psi + \frac{iV_2}{i\hbar} \Psi^* \Psi - \frac{\hbar}{2im} \Psi^* \nabla^2 \Psi + \frac{1}{i\hbar} V_1 \Psi^* \Psi + \frac{iV_2}{i\hbar} \Psi^* \Psi =$$
$$= -\frac{\hbar}{2im} (\Psi^* \nabla^2 \Psi - \Psi \nabla^2 \Psi^*) + \frac{2}{\hbar} V_2 \Psi^* \Psi =$$
$$= -\frac{\hbar}{2im} \nabla \cdot (\Psi^* \nabla \Psi - \Psi \nabla \Psi^*) + \frac{2}{\hbar} V_2 \Psi^* \Psi$$

Θεωρώντας $P = \Psi^* \Psi$ και $\mathbf{J} = \frac{\hbar}{2im} (\Psi^* \nabla \Psi - \Psi \nabla \Psi^*)$ παίρνουμε

$$\frac{\partial P}{\partial t} + \nabla \cdot \mathbf{J} = \frac{2}{\hbar} V_2 P \quad \blacksquare$$

Άσκηση 6.

α) Ισχύει $[l_x, l_y] = i\hbar l_z$, $[l_z, l_x] = i\hbar l_y$, και $[l_z, l_y] = -i\hbar l_x$, άρα

$$\begin{aligned} [l_z, l_+] &= [l_z, l_x + il_y] = [l_z, l_x] + i[l_z, l_y] = \\ &= i\hbar l_y + i(-i\hbar)l_x = \hbar(l_x + il_y) = \hbar l_+ \quad \blacksquare \end{aligned}$$

β)

$$\begin{aligned} l_z \Phi &= l_z l_+ \Psi = l_+ l_z \Psi + \hbar l_+ \Psi = \\ &= l_+ (\lambda \Psi) + \hbar l_+ \Psi = (\lambda + \hbar) l_+ \Psi \end{aligned}$$

$$\begin{aligned} \text{διότι } l_z l_+ - l_+ l_z &= \hbar l_+ \\ \text{και } l_z l_+ &= l_+ l_z + \hbar l_+ \quad \blacksquare \end{aligned}$$

Άσκηση 7.

β)

$$\begin{aligned} \langle \Psi_n, \mathbf{P} \Psi_n \rangle &= \frac{im}{\hbar} \langle \Psi_n [H, \mathbf{r}] \Psi_n \rangle = \\ &= \frac{im}{\hbar} [\langle \Psi_n, H \mathbf{r} \Psi_n \rangle - \langle \Psi_n, \mathbf{r} H \Psi_n \rangle] = \\ &= \frac{im}{\hbar} [E_n \langle \Psi_n, \mathbf{r} \Psi_n \rangle - E_n \langle \Psi_n, \mathbf{r} \Psi_n \rangle] = 0 \quad \blacksquare \end{aligned}$$

γ)

$$\begin{aligned} \langle \Psi_n, H \mathbf{r} \Psi_n \rangle &= \int \Psi_n^* H(\mathbf{r} \Psi_n) d^3x = \\ &= \int (H \Psi_n)^* (\mathbf{r} \Psi_n) d^3x = E_n \int \Psi_n^* (\mathbf{r} \Psi_n) d^3x \quad \blacksquare \end{aligned}$$

Άσκηση 8.

Καταρχήν $F = F(A, x)$ και $[p_x, A] = 0$

Εάν θέλουμε να ορίσουμε την συνάρτηση $g(\hat{B})$ ενός τελεστή \hat{B} θεωρούμε την ανάλυση της $g(x)$ σε σειρά Taylor γύρω από το σημείο $x = 0$:

$$\begin{aligned} g(x) &= \sum_{n=0}^{\infty} \frac{1}{n!} g^{(n)}(x=0) x^n \\ \Rightarrow g(\hat{B}) &= \sum_{n=0}^{\infty} \frac{1}{n!} g^{(n)}(0) \hat{B}^n \end{aligned}$$

Εάν έχουμε την φυσική ποσότητα $C = xp$ τότε ο τελεστής \hat{C} είναι ερμιτιανός εάν ορίσουμε συμμετρικά ως προς τα \hat{x}, \hat{p}

$$\hat{C} = \frac{1}{2} (\hat{x}\hat{p} + \hat{p}\hat{x})$$

Ορίζουμε λοιπόν την ποσότητα $F(A, x)$ ως εξής:

$$F(\hat{A}, \hat{x}) = \frac{1}{2} \left[\sum_{n=0}^{\infty} \frac{1}{n!} \left[F^{(n)}(A, 0) x^n + x^n F^{(n)}(A, 0) \right] \right]$$

όπου

$$F^{(n)}(A, 0) = \left. \frac{d^n F(A, x)}{dx^n} \right|_{x=0}$$

Επομένως

$$\begin{aligned} [p_x, F] &= \left[p_x, \sum_{n=0}^{\infty} \frac{1}{n!} F^{(n)}(A, 0) x^n \right] = \\ &= \sum_{n=0}^{\infty} \frac{1}{n!} \left[[p_x, F^{(n)}(A, 0)] x^n + F^{(n)}(A, 0) [p_x, x^n] \right] = \\ &= 0 + \sum_{n=0}^{\infty} \frac{1}{n!} F^{(n)}(A, 0) (-i\hbar) \frac{\partial x^n}{\partial x} = \\ &= (-i\hbar) \frac{\partial}{\partial x} \left[\sum_{n=0}^{\infty} \frac{1}{n!} F^{(n)}(A, 0) x^n \right] = (-i\hbar) \frac{\partial F}{\partial x} \quad \blacksquare \end{aligned}$$

Άσκηση 9.

Δοκιμάζουμε το ολοκλήρωμα με μια τυχούσα συνάρτηση $f(x)$

$$\begin{aligned} \int_{-\infty}^{+\infty} f(x) \frac{d\zeta}{dx} dx &= \int_{-\infty}^{+\infty} \frac{d}{dx} (f\zeta) dx - \int_{-\infty}^{+\infty} \zeta(x) \frac{df}{dx} dx = \\ &= [f(+\infty)\zeta(+\infty) - f(-\infty)\zeta(-\infty)] - \int_{-\infty}^0 (-a) \frac{df}{dx} dx - \int_0^{+\infty} a \frac{df}{dx} dx = \\ &= [af(+\infty) - (-a)f(-\infty)] + a[f(0) - f(-\infty)] - a[f(+\infty) - f(0)] = \\ &= af(+\infty) - af(+\infty) + af(-\infty) - af(-\infty) + 2af(0) = \\ &= 2af(0) = 2a \int_{-\infty}^{+\infty} f(x) \delta(x) dx \quad \blacksquare \end{aligned}$$

Άσκηση 10.

α)

$$\begin{aligned} e^{-sA} &= \sum_{n=0}^{\infty} \frac{(-sA)^n}{n!} = \sum_{n=0}^{\infty} \frac{(-1)^n s^n}{n!} A^n \\ [e^{-sA}, B] &= \sum_{n=0}^{\infty} \frac{(-1)^n s^n}{n!} [A^n, B] \\ [A^n, B] &= n [A, B] A^{n-1} \\ \Rightarrow [e^{-sA}, B] &= \sum_{n=0}^{\infty} \frac{(-1)^n s^n}{n!} [A, B] n A^{n-1} = \\ &= (-s) [A, B] \sum_{n=1}^{\infty} \frac{(-1)^{n-1} s^{n-1}}{(n-1)!} A^{n-1} = (-s) [A, B] e^{-sA} \quad \blacksquare \end{aligned}$$

β)

$$\begin{aligned} e^{-sA} B e^{sA} &= e^{-sA} B e^{sA} + B e^{-sA} e^{sA} - B e^{-sA} e^{sA} = \\ &= [e^{-sA}, B] e^{sA} + B e^{-sA} e^{sA} = \\ &= -s [A, B] e^{-sA} e^{sA} + B e^{-sA} e^{sA} = -s [A, B] + B \end{aligned}$$

$$\begin{aligned} e^{-sA} e^{sA} &= \left[1 - sA + \frac{(sA)^2}{2!} - \frac{(sA)^3}{3!} + \dots \right] \left[1 + sA + \frac{(sA)^2}{2!} + \frac{(sA)^3}{3!} + \dots \right] = \\ &= 1 + (sA - sA) + \left[\frac{(sA)^2}{2} - (sA)^2 + \frac{(sA)^2}{2} \right] + s^3 [\dots] + \dots = \\ &= 1 + 0 + 0 + 0 + \dots = 1 \quad \blacksquare \end{aligned}$$

Άσκηση 11.

α)

$$A = A(\omega)$$
$$\frac{dA}{d\omega} = \lim_{\Delta\omega \rightarrow 0} \frac{A(\omega + \Delta\omega) - A(\omega)}{\Delta\omega} \quad \blacksquare$$

β)

$$\begin{aligned} \frac{d(AB)}{d\omega} &= \lim_{\Delta\omega \rightarrow 0} \frac{A(\omega + \Delta\omega)B(\omega + \Delta\omega) - A(\omega)B(\omega)}{\Delta\omega} = \\ &= \lim_{\Delta\omega \rightarrow 0} \frac{1}{\Delta\omega} [A(\omega + \Delta\omega)B(\omega + \Delta\omega) + A(\omega)B(\omega + \Delta\omega) - A(\omega)B(\omega + \Delta\omega) - A(\omega)B(\omega)] = \\ &= \lim_{\Delta\omega \rightarrow 0} \frac{1}{\Delta\omega} [A(\omega + \Delta\omega) - A(\omega)] B(\omega + \Delta\omega) + \lim_{\Delta\omega \rightarrow 0} \frac{A(\omega)}{\Delta\omega} [B(\omega + \Delta\omega) - B(\omega)] = \\ &= \frac{dA}{d\omega} B + A \frac{dB}{d\omega} \quad \blacksquare \end{aligned}$$

γ)

$$\begin{aligned} AA^{-1} = 1 &\Rightarrow \frac{d}{d\omega}(AA^{-1}) = 0 \\ \Rightarrow \frac{dA}{d\omega} A^{-1} + A \frac{dA^{-1}}{d\omega} &= 0 \\ \Rightarrow \frac{dA}{d\omega} A^{-1} = -A \frac{dA^{-1}}{d\omega} &\Rightarrow \frac{dA^{-1}}{d\omega} = -A^{-1} \frac{dA}{d\omega} A^{-1} \\ \Rightarrow \frac{dA}{d\omega} = -A \frac{dA^{-1}}{d\omega} A &\quad \blacksquare \end{aligned}$$